

Supporting Information

Large enhancement of magnetocaloric and barocaloric effect by hydrostatic pressure in

La(Fe_{0.92}Co_{0.08})_{11.9}Si_{1.1} with NaZn₁₃-type structure

Jiazheng Hao, Fengxia Hu,* Jian-Tao Wang, Fei-Ran Shen, Zibing Yu, Houbo Zhou, Hui Wu, Qingzhen Huang, Kaiming Qiao, Jing Wang,* Jun He, Lunhua He, Ji-Rong Sun, and Baogen Shen*

Series title:

- I. The neutron powder diffraction (NPD) pattern collected at 300K in the absence of pressure for La(Fe_{0.92}Co_{0.08})_{11.9}Si_{1.1} compound.**
- II. Crystal structure parameters of La(Fe_{0.92}Co_{0.08})_{11.9}Si_{1.1} compound under different pressures obtained from the refinements of power neutron diffraction data.**
- III. The total and partial electronic DOS of up spins and down spins for La(Fe_{0.92}Co_{0.08})_{11.9}Si_{1.1} compound in the ferromagnetic state at ambient pressure.**
- IV. Evaluation of the enlarged $\Delta V/V$ to the enhancement of BCE by hydrostatic pressure.**

I. The neutron powder diffraction (NPD) pattern collected at 300K in the absence of pressure for $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ compound.

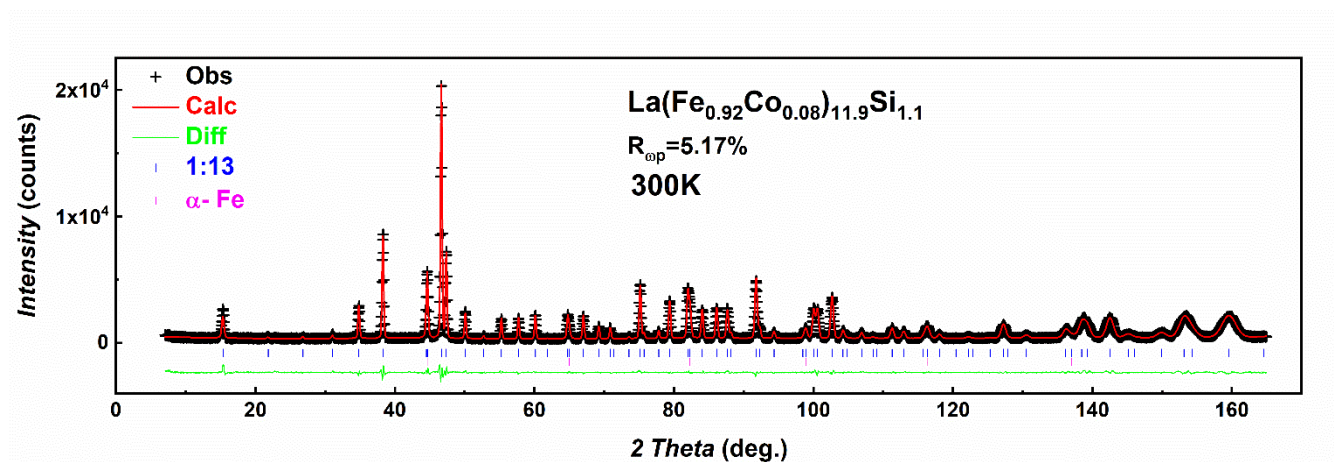


Figure S1. The collected NPD pattern together with the Rietveld refinement plot at 300K in the absence of pressure for $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ compound. The black cross and red line represent observed and calculated profiles, respectively. The vertical bars indicate the Bragg reflection positions of NaZn_{13} -phase (blue) and α -Fe (purple).

II. Crystal structure parameters of $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ compound under different pressures

obtained from the refinements of power neutron diffraction data.

Table S1. Crystal structure parameters of ferromagnetic (FM) and paramagnetic (PM) phase of $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ compound under different pressures obtained from the refinements of power neutron diffraction data: the lattice constant a , unit cell volume V , variation of lattice parameters and unit cell volumes across phase transition, the fractional coordinates at 96i site y , z , La-Fe₂ atomic distance, intra-icosahedral Fe-Fe distances (B1, B2, B3) and inter-icosahedral ones (B4, B5), angles of Fe-Fe-Fe atomic chain (intra-icosahedron: Ang-1', Ang-2', Ang-3', Ang-4', on surface of the icosahedron (Ang-1, Ang-2), inter-icosahedron (Ang-3, Ang-4, Ang-5, Ang-6, Ang-7), **Figure 4b**), and the reliability factors R_{wp} .

| | 0kbar | | | 3kbar | | | 6kbar | | | 9kbar | | |
|-------------------------------------|--------------|--------------|-------------|--------------|--------------|-------------|--------------|--------------|-------------|--------------|--------------|-------------|
| | FM (242K) | PM (338K) | Diff (%) | FM (201K) | PM (314K) | Diff (%) | FM (176K) | PM (302K) | Diff (%) | FM (149K) | PM (284K) | Diff (%) |
| $a(\text{\AA})$ | 11.5333 (3) | 11.4908 (3) | -0.37 | 11.5237 (3) | 11.4745 (3) | -0.43 | 11.5132 (3) | 11.4564 (3) | -0.50 | 11.5052 (3) | 11.4457 (4) | -0.52 |
| $V(\text{\AA}^3)$ | 1533.6 (2) | 1517.2 (1) | -1.08 | 1530.3 (1) | 1510.8 (1) | -1.29 | 1526.1 (1) | 1503.6 (1) | -1.50 | 1522.9 (1) | 1499.4 (1) | -1.57 |
| y/a | 0.1823 (3) | 0.1806 (3) | -0.94 | 0.1822 (3) | 0.1801 (3) | -1.22 | 0.1816 (3) | 0.1793 (3) | -1.28 | 0.1809 (3) | 0.1794 (3) | -0.82 |
| z/a | 0.1198 (3) | 0.1193 (3) | -0.45 | 0.1198 (3) | 0.1195 (3) | -0.27 | 0.1193 (3) | 0.1186 (3) | -0.60 | 0.1201 (3) | 0.1187 (3) | -1.15 |
| La-Fe ₂ (\AA) | 3.3431(18) | 3.3382(18) | -0.15 | 3.3406(18) | 3.3339(18) | -0.2 | 3.342(18) | 3.3355(18) | -0.19 | 3.3375(18) | 3.3276(18) | -0.30 |
| B1(\AA) | 2.517(3) | 2.487(3) | -1.21 | 2.514(3) | 2.48(3) | -1.37 | 2.502(3) | 2.463(3) | -1.58 | 2.501(3) | 2.459(3) | -1.71 |
| B2(\AA) | 2.775(5) | 2.742(5) | -1.20 | 2.766(5) | 2.729(5) | -1.36 | 2.758(5) | 2.721(5) | -1.36 | 2.755(5) | 2.712(5) | -1.59 |
| B3(\AA) | 2.618(3) | 2.585(3) | -1.28 | 2.615(3) | 2.576(3) | -1.51 | 2.602(3) | 2.563(3) | -1.52 | 2.599(3) | 2.556(3) | -1.68 |
| B4(\AA) | 2.384(5) | 2.405(5) | 0.87 | 2.392(5) | 2.407(5) | 0.62 | 2.393(5) | 2.415(5) | 0.91 | 2.394(5) | 2.413(5) | 0.79 |
| B5(\AA) | 2.500(4) | 2.510(4) | 0.40 | 2.499(4) | 2.517(4) | 0.72 | 2.504(4) | 2.513(4) | 0.36 | 2.511(4) | 2.519(4) | 0.32 |
| Ang-1'($^\circ$) | 66.72(17) | 67.16(17) | 0.66 | 66.62(17) | 66.69(17) | 0.11 | 66.60(17) | 66.97(17) | 0.56 | 67.17(17) | 66.98(17) | -0.28 |
| Ang-2' ($^\circ$) | 56.66(15) | 56.42(15) | 0.42 | 56.69(15) | 56.66(15) | -0.05 | 56.70(15) | 56.52(15) | -0.32 | 56.42(15) | 56.51(15) | 0.16 |
| Ang-3' ($^\circ$) | 62.66(11) | 62.56(11) | -0.16 | 62.68(11) | 62.67(11) | -0.02 | 62.68(11) | 62.60(11) | -0.13 | 62.56(11) | 62.60(11) | 0.06 |
| Ang-4' ($^\circ$) | 58.69(11) | 58.72(11) | 0.05 | 58.66(11) | 58.67(11) | 0.02 | 58.66(11) | 58.70(11) | 0.07 | 58.72(11) | 58.70(11) | -0.03 |
| Ang-1 ($^\circ$) | 58.07(13) | 57.82(13) | -0.43 | 58.13(13) | 58.09(13) | -0.07 | 58.14(13) | 57.93(13) | -0.36 | 57.82(13) | 57.92(13) | 0.17 |
| Ang-2 ($^\circ$) | 63.85(15) | 64.36(15) | 0.80 | 63.73(15) | 63.81(15) | 0.13 | 63.71(15) | 64.14(15) | 0.67 | 64.37(15) | 64.15(15) | -0.34 |
| Ang-3 ($^\circ$) | 56.48(16) | 56.92(16) | 0.78 | 56.47(16) | 56.93(16) | 0.81 | 56.68(16) | 57.27(16) | 1.04 | 56.74(16) | 57.24(16) | 0.88 |
| Ang-4 ($^\circ$) | 67.03(18) | 66.17(18) | -1.28 | 67.06(18) | 66.14(18) | -1.37 | 66.63(18) | 65.45(18) | -1.77 | 66.53(18) | 65.52(18) | -1.52 |
| Ang-5 ($^\circ$) | 55.66(15) | 56.15(15) | 0.88 | 55.71(15) | 56.50(15) | 1.42 | 56.11(15) | 56.92(15) | 1.44 | 55.81(15) | 56.85(15) | 1.86 |
| Ang-6 ($^\circ$) | 64.52(16) | 63.25(16) | -1.97 | 64.60(16) | 63.45(16) | -1.78 | 64.08(16) | 62.50(16) | -2.47 | 63.68(16) | 62.57(16) | -1.74 |
| Ang-7 ($^\circ$) | 59.82(16) | 60.60(16) | 1.30 | 59.70(16) | 60.05(16) | 0.59 | 59.8(16) | 60.58(16) | 1.30 | 60.51(16) | 60.58(16) | 0.12 |
| $R_{\text{wp}}(\%)$ | 6.78 | 6.44 | - | 7.29 | 6.49 | - | 6.98 | 5.99 | - | 6.63 | 6.85 | - |

III. The total and partial electronic DOS of up spins and down spins for $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ compound in the ferromagnetic state at ambient pressure.

It is worth noting that a very denser energy point is required to describe the DOSs near the Fermi level for such large and complicated magnetic compound. All of the DOSs reported in this work are plotted with an energy-range from -8 eV to 4 eV with 12000 energy points (NEDOS=12000) and a small broadening energy of 0.01 eV (SIGMA= 0.01).

To understand the mechanism of magnetic behavior of $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$, we have calculated the total and partial electronic density of states (DOS), as shown in **Figure S2**. The Fermi energy is shifted to zero ($E_F=0$). There is a large splitting between the spin-up and spin-down DOS below the E_F , as shown in **Figure S2a**, which result a large magnetic moment of 210.16 μ_B at ambient pressure. The spin-up bands is almost completely filled, while the spin-down bands is partially filled with some spin-down holes, thus using the familiar terminology, we can consider the $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ compound as strong ferromagnetism. The spin-up and spin-down DOS peaks from -5 eV to 0 eV are mainly from the 3d bands of Fe and Co atoms (see **Figure S2c** and **Figure S2d**), and the splittings result a local magnetic moment of 1.91~2.56 μ_B on the Fe sites and 1.41~1.65 μ_B on the Co sites, respectively. There is a small induced splitting of 5d bands of La around the energy of -1.5 eV (which induce a small magnetic moment of -0.21 μ_B on La sites), indicating the La-5d and (Fe,Co)-3d bands are well hybridized, but the main part of the partial DOS of the La-5d is located above the Fermi level (**Figure S2b**). Also there is a small induced splitting of 2p bands of Si around the energy range from -6.0 eV to -1.0 eV, as shown in **Figure S2e**, which induce a small magnetic moment of -0.11 μ_B on Si sites. Meanwhile the peaks from the 3s bands of Si are located around the energy of -8 eV (**Figure S2e**), lying far below E_F and the peaks from the 4f bands of La are around the energy of 3 eV, above E_F (**Figure S2b**).

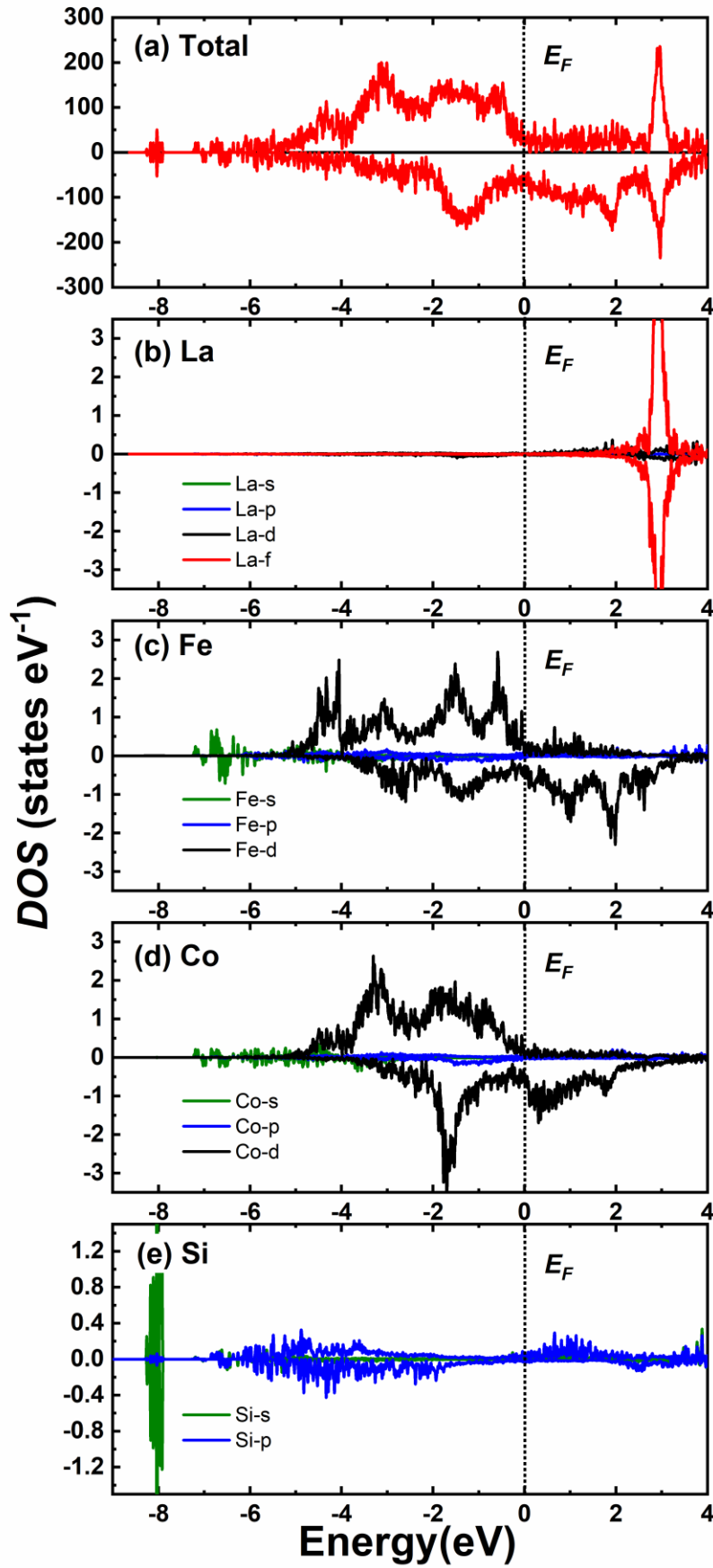


Figure S2. The total and partial electronic DOS of up spins and down spins for $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ compound in the ferromagnetic state at ambient pressure. The Fermi energy is shifted to zero.

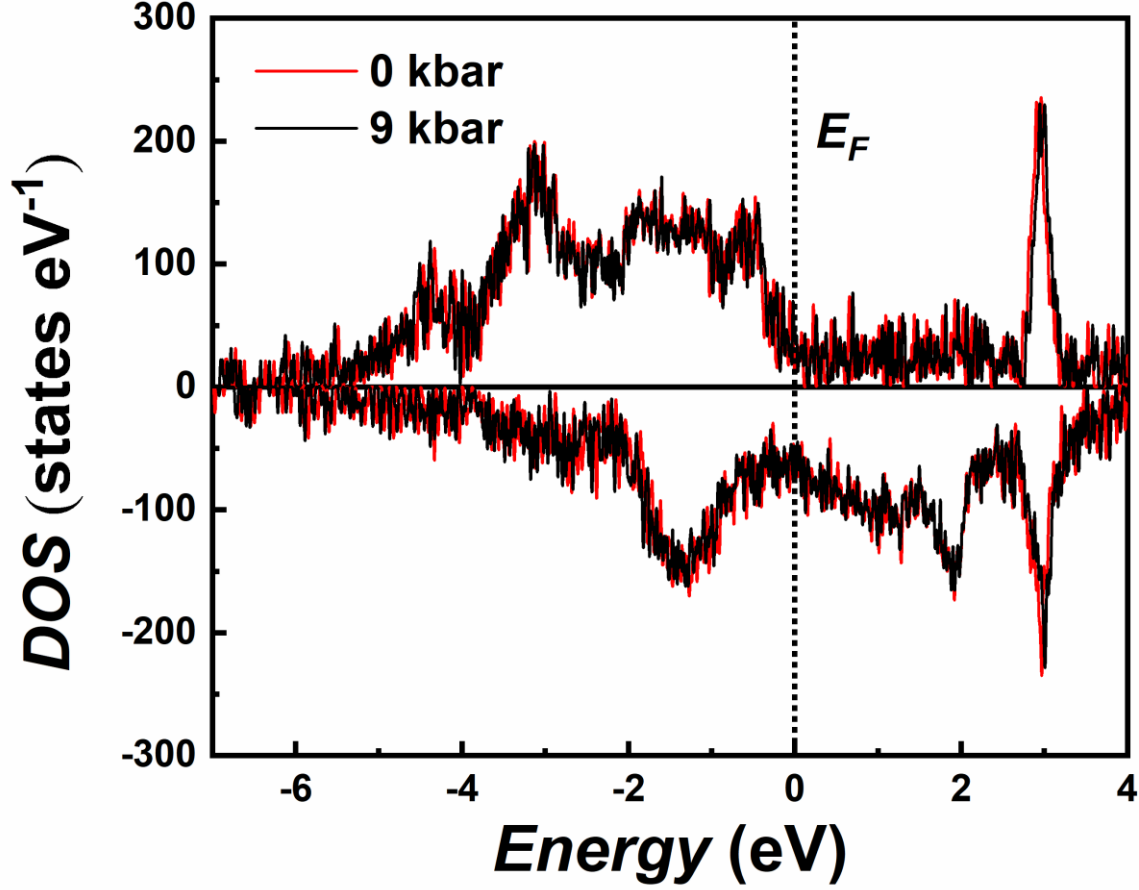


Figure S3. The total density of states of $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ compound in the ferromagnetic state under pressures 0 and 9 kbar. The Fermi energy is shifted to zero.

IV. Evaluation of the enlarged $\Delta V/V$ to the enhancement of BCE by hydrostatic pressure.

The barocaloric effect (BCE) of $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ shows a rapid increase with pressure due to the enhanced magnetovolume effect and the evolution of the transition nature by hydrostatic pressure. To understand the essence, we quantitatively evaluate the contribution of enlarged volume change $\Delta V/V$ during phase transition to the enhancement of BCE. Firstly, we assume that the effect of pressure on the movement and evolution of phase transition order remains unchanged but does not affect the $\Delta V/V$. Then the lattice entropy $S'_{\text{Latt-PT}}$ involving phase transition under different pressures can be calculated based on Debye approximation, as shown in **Figure S4a**. By comparing the $S'_{\text{Latt-PT}}$ with and without pressure, we can get the pressure-induced lattice entropy change $\Delta S'_{\text{Latt}}$, i.e. the contribution of lattice entropy to the barocaloric effect ($\Delta S'_{\text{BCE-Latt}}$), as plotted in **Figure S4b**. We also suppose that the spin contribution to the

$S_{\text{Tot-PT}}$ remains nearly independent of pressure (**Figure 7d**). Then total entropy $S'_{\text{Tot-PT}}$ involving phase transition (**Figure S4c**), as well as the pressure-induced total entropy change $\Delta S'_{\text{BCE-Tot}}$ (**Figure S4d**), i.e. the barocaloric effect (BCE') without considering the effect of enlarged $\Delta V/V$, can be obtained. **Figure 7g** and **Figure 7h** depict the comparison between $\Delta S'_{\text{BCE-Latt}}$ and $\Delta S_{\text{BCE-Latt}}$, and between $\Delta S'_{\text{BCE-Tot}}$ and $\Delta S_{\text{BCE-Tot}}$, respectively, where the shaded area denotes the contribution from enlarged $\Delta V/V$ under 9kbar. One can note that, compared to $\Delta S'_{\text{BCE-Latt}}$ and $\Delta S'_{\text{BCE-Tot}}$, the maximum of $\Delta S_{\text{BCE-Latt}}$ and $\Delta S_{\text{BCE-Tot}}$ curves increases by 39% and 22%, while the area under the curves which connects RC enlarges by 47% and 20%, respectively, due to the contribution of enlarged $\Delta V/V$ by hydrostatic pressure.

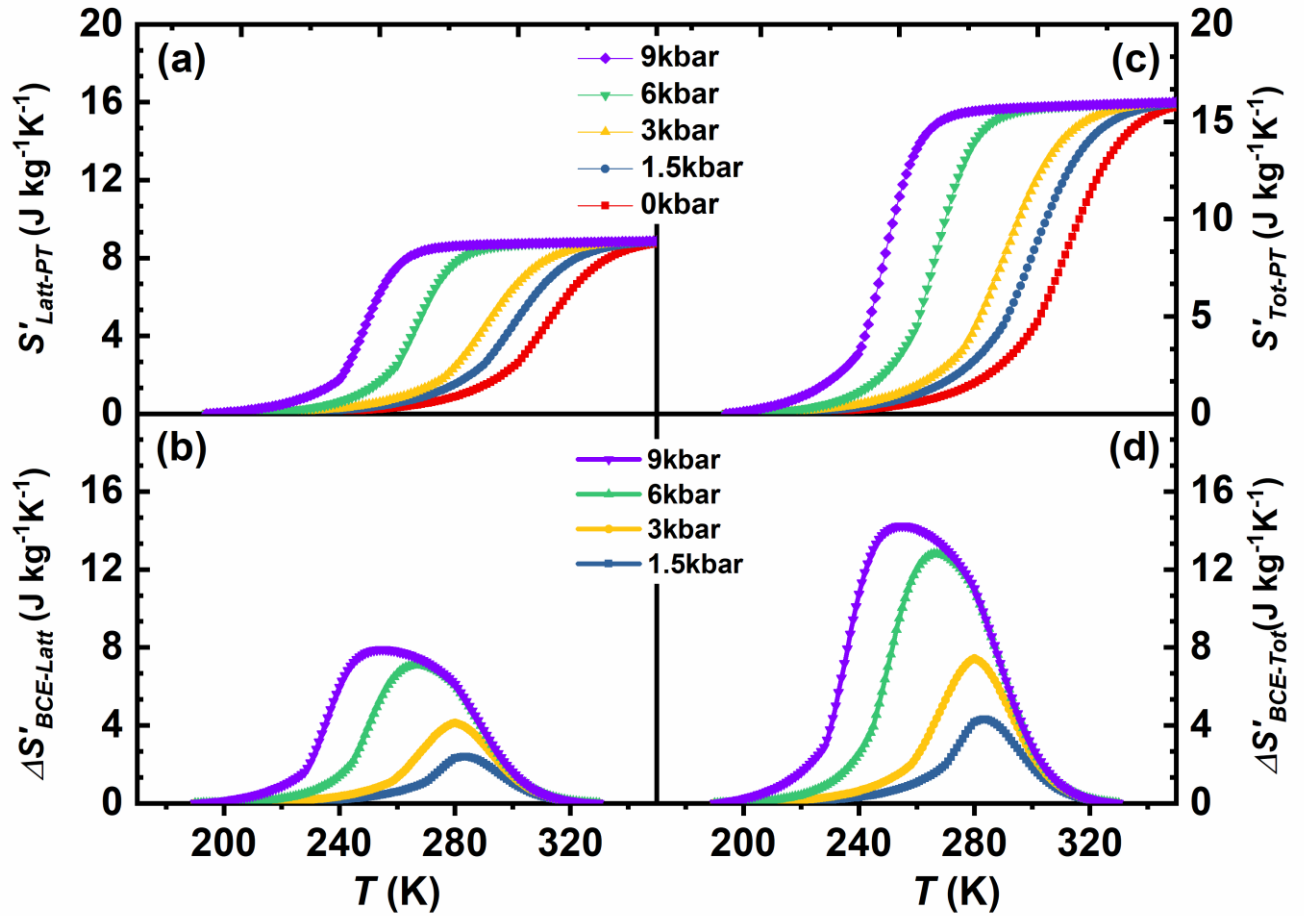


Figure S4. Supposing that the effect of pressure on the movement and evolution of phase transition order remains unchanged but does not affect the $\Delta V/V$, (a) the calculated lattice entropy $S'_{\text{Latt-PT}}$ involving phase transition based on Debye approximation, (b) pressure-induced lattice entropy change $\Delta S'_{\text{BCE-Latt}}$, i.e. the lattice contribution to barocaloric effect ($\Delta S'_{\text{BCE-Latt}}$), (c) the total entropy $S'_{\text{Tot-PT}}$ involving phase transition only under different pressures, and (d) the pressure-induced total entropy change $\Delta S'_{\text{BCE-Tot}}$, i.e. the barocaloric effect (BCE') without considering the effect of enlarged $\Delta V/V$.

IV. Debye temperature Θ_D of $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ in the absence of pressure.

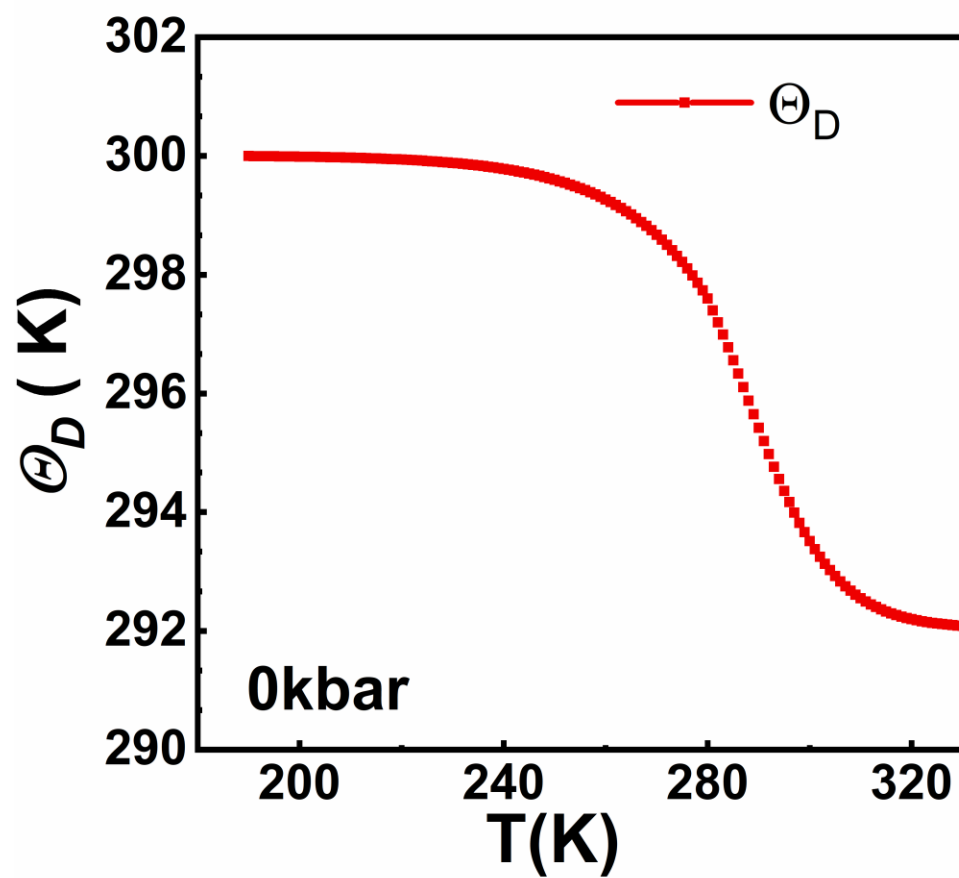


Figure S5 Temperature dependence of Θ_D for $\text{La}(\text{Fe}_{0.92}\text{Co}_{0.08})_{11.9}\text{Si}_{1.1}$ calculated by Debye approximation in the absence of pressure.